

## Research Article

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# Study on the formation mechanism of the hard-shell layer of liquefied silty soil

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**Abstract:** Hard-shell layers often form during the migration and segregation of fine-grained sediments and play a critical role in influencing the mechanical behavior of seabed soils, with implications for marine geological hazards. Understanding the mechanisms that govern the formation and evolution of these layers is therefore essential for coastal and estuarine geotechnical engineering. In this study, laboratory experiments were conducted to investigate the formation mechanisms of hard-shell layers in saturated silty soils subjected to two types of controlled physical disturbances: vertical hammering and horizontal rotating wheel-induced shear. Using high-speed particle image velocimetry, pore pressure sensors, and earth pressure measurements, the study analyzed the coupled evolution of the force field, particle migration field, and liquefaction behavior under different disturbance modes. Results show that liquefaction is a prerequisite for hard-shell layer formation, allowing fine particles to migrate upward while coarse particles settle, leading to effective particle separation. Hammering produced high pore pressure and strong vertical particle mobility,

resulting in a thin but dense hard-shell layer with high strength. In contrast, the rotating wheel disturbance generated moderate liquefaction and vortex-driven particle segregation, forming a thicker hard-shell layer with relatively lower mechanical strength. The findings highlight a trade-off between shell thickness and strength depending on the disturbance type and suggest that the mechanism of fine-coarse particle separation under liquefaction conditions is key to shell formation. This research provides a theoretical and experimental basis for understanding internal liquefaction processes in silty coastal slopes, particularly under the influence of waves or seismic activity following landslides.

**Keywords:** liquefaction-induced sedimentation, hard-shell layer, physical disturbance, formation mechanism

## 1 Introduction

The shallow marine strata along silty shores are primarily composed of fine-grained silt and clay particles with small particle sizes and narrow grading distributions [1]. In coastal zones, extensive land reclamation projects have utilized locally dredged seabed silt as fill material [2–4]. However, due to the high clay content that remains suspended during the hydraulic fill process, the resulting sediment layer often exhibits a prolonged consolidation period, low shear strength, and poor stability [5]. Furthermore, the mechanical behavior of silty soils is highly sensitive to cyclic loading conditions, such as those induced by waves and tides. Studies have shown that fine-grained content, particularly clay, significantly influences the pore water pressure response and shear strength degradation under dynamic excitation [6,7].

According to classical soil mechanics and liquefaction theory, cyclic loading in saturated loose granular soils can lead to the build-up of excess pore water pressure and a corresponding loss of effective stress, ultimately resulting in a liquefied state. This phenomenon is particularly pronounced in silty soils due to their dual characteristics: low permeability impedes pore pressure dissipation, while low plasticity limits structural resistance. During liquefaction,

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suspended clay particles migrate under the influence of seepage forces and buoyancy, facilitating internal segregation and the formation of preferential flow channels within the soil skeleton [8]. Upon cessation of loading, pore pressure gradually dissipates, and the soil structure undergoes thixotropic recovery, sometimes leading to the formation of denser, more stable layers within the previously liquefied stratum [9–11]. These densified layers can potentially serve as natural or engineered rigid crusts if properly controlled.

However, existing studies lack a mechanistic understanding of how artificial liquefaction can be harnessed to reliably form high-strength shell layers within soft silt foundations, especially under coupled wave and vibration disturbances. While quantitative research has explored artificial liquefaction [12], its engineering application for *in situ* soil improvement remains unclear due to insufficient theoretical integration and limited laboratory validation.

This study aims to investigate whether artificial disturbance – induced through vertical vibration and horizontal rolling – can accelerate the formation of a dense, rigid silt layer and enhance the strength and permeability of liquefied sediments.

Two primary liquefaction mechanisms are considered: (1) vertical vibration-induced liquefaction simulating seismic activity and (2) horizontal cyclic shear induced by wave motion [13,14]. In the vertical vibration scenario, variables such as clay content, dry density, consolidation stress, and dynamic stress ratio influence the cyclic liquefaction resistance of the soil [2,6,7]. Meanwhile, in wave-induced liquefaction, flume experiments reveal that repeated shear from wave forces can redistribute particles, increasing density and reducing pore volumes [15]. Both types of disturbances are rooted in fundamental principles of cyclic mobility and stress path theory in saturated soils, emphasizing the role of drainage conditions, grain-size distribution, and fabric structure.

To explore the interaction between these mechanisms, this study employs high-speed particle image velocimetry (PIV) in conjunction with pore pressure and earth pressure sensors, enabling detailed observation of particle migration and pressure response. A 1.2 m water tank simulates controlled disturbance environments to test both isolated and combined liquefaction modes. This setup allows real-time tracking of the transformation from loose to dense soil structures, shedding light on the post-liquefaction densification process and the potential formation of rigid strata under cyclic wave loading.

By elucidating the interplay between artificial liquefaction, pore pressure dissipation, and soil structure adjustment, this research aims to provide a novel soil improvement

method for estuarine and coastal engineering projects. The findings offer a theoretical and practical basis for reducing the consolidation time, increasing permeability, and lowering the vacuum preloading cost in soft silt foundations. Furthermore, the study may enhance our understanding of post-landslide liquefaction dynamics in silty slopes adjacent to coastlines subjected to seismic or wave disturbances.

## 2 Materials and methods

### 2.1 Soil materials

The soil materials in this experiment were silty soil samples from the Yellow River Delta region. The measurement showed a water content of about 30%, with a  $d_{50}$  of 0.038 mm. The grading ratio is shown in Figure 1.

### 2.2 PIV flume experiment

#### 2.2.1 PIV test procedure

The PIV flume experiment was conducted to assess the force field and flow field structure of the liquefied soil, as illustrated in Figure 2. The silty soil layer deposited in a  $0.6\text{ m} \times 0.6\text{ m} \times 0.6\text{ m}$  flume was disturbed by a 79 g hammer body (hammer height of 20 cm) and a submerged motor with a 10 mm diameter rotating wheel (20 rpm). Mechanical indices related to soil hydration were measured using vertical and horizontal pore pressure sensors and soil pressure sensors embedded at the bottom and right center of the flume. Additionally, PIV technology was applied to measure the two-dimensional flow field.

The light source for the experiment was an intense light unit with a power of 500 W, located 1 m in front of the flume. Fine sandstones were used as tracer particles. The CCD camera was located 0.5 m from the observation surface of the flume and has a resolution of  $1,600 \times 1,200$  pixels and a frame rate of 16 fps. The test area of  $30\text{ cm} \times 25\text{ cm}$  was in the middle of the flume.

The migration velocity of the tracer particles was determined by comparing two photos before and after each frame using TSI-Insight3G PIV post-processing software. Particles with the same velocity were connected to form the velocity gradient. The detailed interference pattern is displayed in Figure 2. After adding soil and leaving it

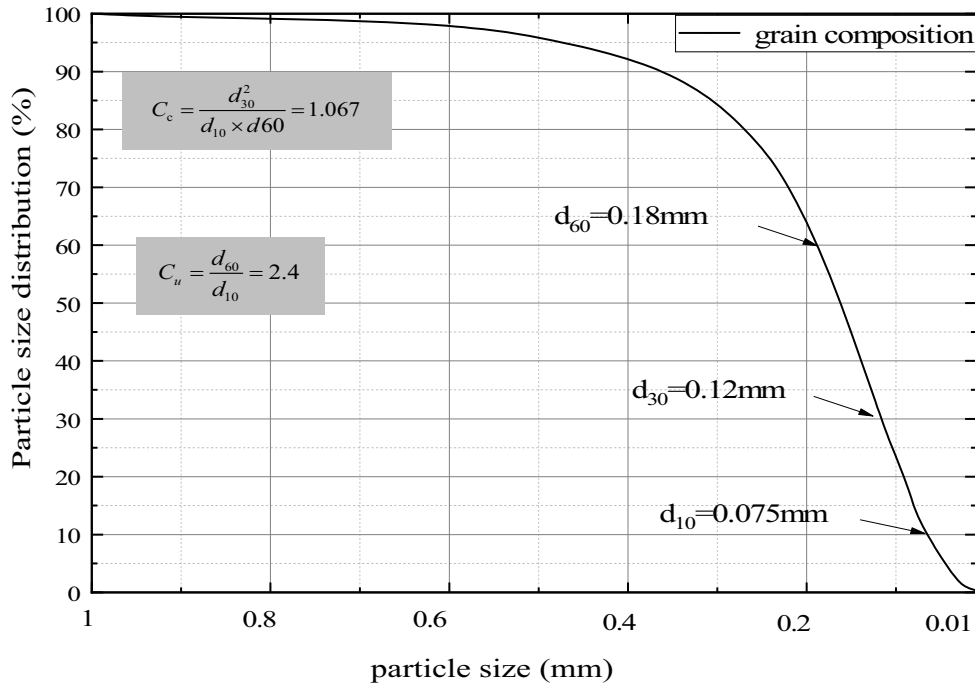


Figure 1: Initial grading.

for 12 h, the sedimentary strata were hammered at 10 times/min and 10 rpm for the PIV hammer test and the PIV disturbance experiment, respectively.

## 2.2.2 Detailed implementation of PIV post-processing and data extraction

The post-processing of PIV data was conducted using PIVlab (version X.X), a MATLAB-based open-source software for 2D PIV analysis. A multi-pass fast Fourier transform cross-correlation algorithm was used to calculate the displacement fields between consecutive image pairs.

### 2.2.2.1 Image preprocessing and correlation analysis

To enhance image contrast and minimize noise, histogram equalization and Gaussian filtering were applied prior to correlation analysis. The PIV processing followed a three-step iterative scheme with decreasing interrogation window sizes ( $64 \times 64 \rightarrow 32 \times 32 \rightarrow 16 \times 16$  pixels) and 50% overlap. The displacement vector  $\vec{u} = (u_x, u_y)$  at each point was determined by finding the peak of the cross-correlation function:

$$R(\Delta x, \Delta y) = \sum_{i,j} I_1(i, j) \cdot I_2(i + \Delta x, j + \Delta y), \quad (1)$$

where  $I_1$  and  $I_2$  are the grayscale intensity values of the reference and deformed images, respectively.

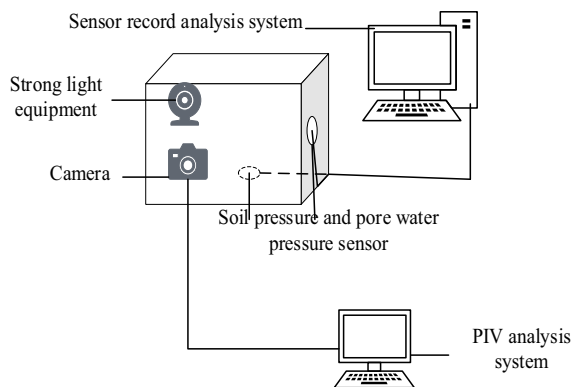
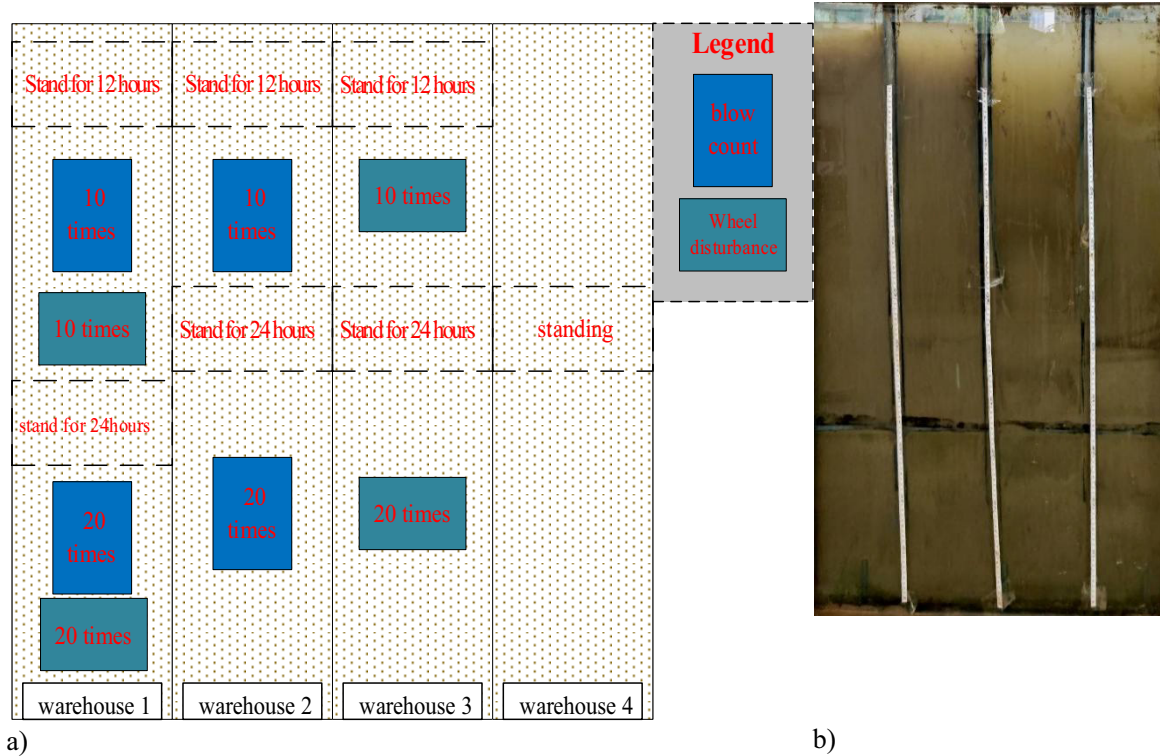


Figure 2: Schematic diagram of the PIV flume experiment.



**Figure 3:** Schematic diagram of the 1.2 m flume disturbance scheme: (a) schematic diagram of the flume and (b) flume for indoor flume experiments.

Sub-pixel accuracy was achieved using Gaussian peak fitting around the cross-correlation maximum. Outlier vectors were filtered using the normalized median test.

### 2.2.2.2 Displacement-to-strain field transformation

The displacement fields were converted to strain fields using finite difference approximations. The normal and shear strain components were computed as follows:

normal strains

$$\varepsilon_{xx} = \frac{\partial u_x}{\partial x}, \quad \varepsilon_{yy} = \frac{\partial u_y}{\partial y}, \quad (2)$$

shear strain:

$$\gamma_{xy} = \frac{\partial u_x}{\partial y} + \frac{\partial u_y}{\partial x}. \quad (3)$$

The derivatives were estimated using a central difference scheme on a uniform grid.

### 2.2.2.3 Principal strain analysis

From the computed strain tensor, the principal strains  $\varepsilon_1$  and  $\varepsilon_2$  were extracted using eigenvalue analysis:

$$\varepsilon_{1,2} = \frac{\varepsilon_{xx} + \varepsilon_{yy}}{2} \pm \sqrt{\left(\frac{\varepsilon_{xx} - \varepsilon_{yy}}{2}\right)^2 + \left(\frac{\gamma_{xy}}{2}\right)^2}. \quad (4)$$

This allows the identification of zones of maximum extension and compression, which are critical for understanding crack initiation and propagation.

### 2.2.2.4 Data extraction and visualization

For each loading increment, the displacement and strain fields were exported as matrices from PIVlab. The evolution of deformation was tracked by extracting:

- maximum principal strain values at each frame,
- displacement profiles along specific lines of interest (e.g., specimen mid-height),
- time-dependent strain localization zones, and
- crack-tip propagation paths using strain threshold (typically where  $\varepsilon_1 > 0.5\%$ ).

These datasets were further processed using MATLAB scripts to generate contour maps, vector fields, and strain evolution plots, enabling quantitative comparison between different loading stages.

## 2.3 Comparison test of the 1.2 m high flume

Due to the small size of the PIV test flume, it was impossible to obtain enough samples to comprehensively analyze the distribution of soil particles and changes in physical properties after the disturbance. Therefore, a separation

experiment was conducted in a 1.2 m high flume, as shown in Figure 3. According to the PIV test, the disturbed area was only 5 cm × 5 cm. Since the disturbance range of the soil was five times the diameter of the disturbance wheel, a 50 cm diameter disturbance wheel was used in the 1.2 m high flume. The flume was divided into four experimental compartments with clapboards (Warehouse 1, Warehouse 2, Warehouse 3, and Warehouse 4), each measuring 0.5 m × 0.5 m × 1.2 m.

In Warehouse 1, the stratum was hammered and rolled 10 times 12 h after soil addition and continuously disturbed with 20 hammer strikes and 20 rolls 24 h after soil addition. In Warehouse 2, the stratum was hammered 10 times and 20 times 12 and 24 h after soil addition, respectively. In Warehouse 3, the stratum was rolled 10 times and 20 times 10 and 24 h after soil addition, respectively. Warehouse 4 was used for static observation.

During the flume experiment, a 10 kg disturbance hammer and a 10 cm diameter grid wheel were used. The disturbance modes were joint disturbance for Warehouse 1 (hammer and wheel), hammer-promoted settlement for Warehouse 2, wheel-promoted settlement for Warehouse 3, and static state for Warehouse 4. The specific disturbance schemes are illustrated in Figure 3.

The bearing capacity of the foundation at different layers was measured at 24 and 48 h of soil settlement after the disturbance treatment. Soil samples were collected at the formation of the hard-shell layer and tested for physical parameter indices, real-time bearing capacity, and rheological property indices.

Figure 3 presents the schematic and physical setup of the 1.2 m flume used in indoor disturbance experiments. Subfigure (a) illustrates the layout and zoning of the flume system, including key dimensions and disturbance areas. Subfigure (b) shows the actual flume used in testing, designed to simulate field-like disturbance conditions. This figure provides essential context for understanding the controlled experimental environment and ensures that readers can interpret the test conditions and setup independently from the main text, thereby enhancing scientific clarity and rigor.

### 3 Analysis of experimental results

Through physical disturbance tests on remolded silty soil samples extracted from the Yellow River Delta, flow field structures associated with specific disturbance modes were revealed. These disturbances significantly increase the pore water pressure within the soil samples, creating a cascade of dynamic changes. Data were obtained through sensors in the

experimental setup. Sensors positioned at the base of the samples recorded the interaction between the bottom soil pressure ( $P_{\text{bottom}}$ ) and the corresponding bottom pore water pressure ( $u_{\text{bottom}}$ ). Meanwhile, sensors on the lateral walls recorded the lateral soil pressure ( $P_{\text{side}}$ ) and its associated lateral pore water pressure ( $u_{\text{side}}$ ). These findings are consistent with existing studies that the pore water pressure in silty soil increases with the effect of dynamic loading. In turn, effective stress gradually decreases until a critical point is reached, leading to liquefaction of the soil samples. Notably, the horizontal pore water pressure consistently exceeds the vertical pore water pressure in the experiments. Additionally, the vertical soil pressure is consistently lower than the corresponding vertical pore soil pressure. These differences demonstrate the interactions between saturation, hydration, and structural response within the soil samples. In summary, the behavior of remolded silty soil under dynamic loads obtained through tests and analysis elucidates the interactions between pore water pressure, effective stress, and the onset of soil liquefaction, facilitating the understanding of soil mechanics in response to various disturbance regimes.

#### 3.1 Analysis of interference force field

The effective soil pressure  $P$  and the pore water pressure  $u$  show a certain periodicity with the dynamic load during the PIV flume disturbance. Therefore, the bottom soil pressure  $P_{\text{bottom}}$ , side soil pressure  $P_{\text{side}}$ , bottom pore water pressure  $u_{\text{bottom}}$ , and side pore water pressure  $u_{\text{side}}$  were evaluated under hammering and wheel conditions. In this study, the liquefaction disturbance of mixed silty soil was recorded, and the force field changes of the liquefaction disturbance of silty soils were measured step-by-step using a pore pressure sensor and a soil pressure sensor.

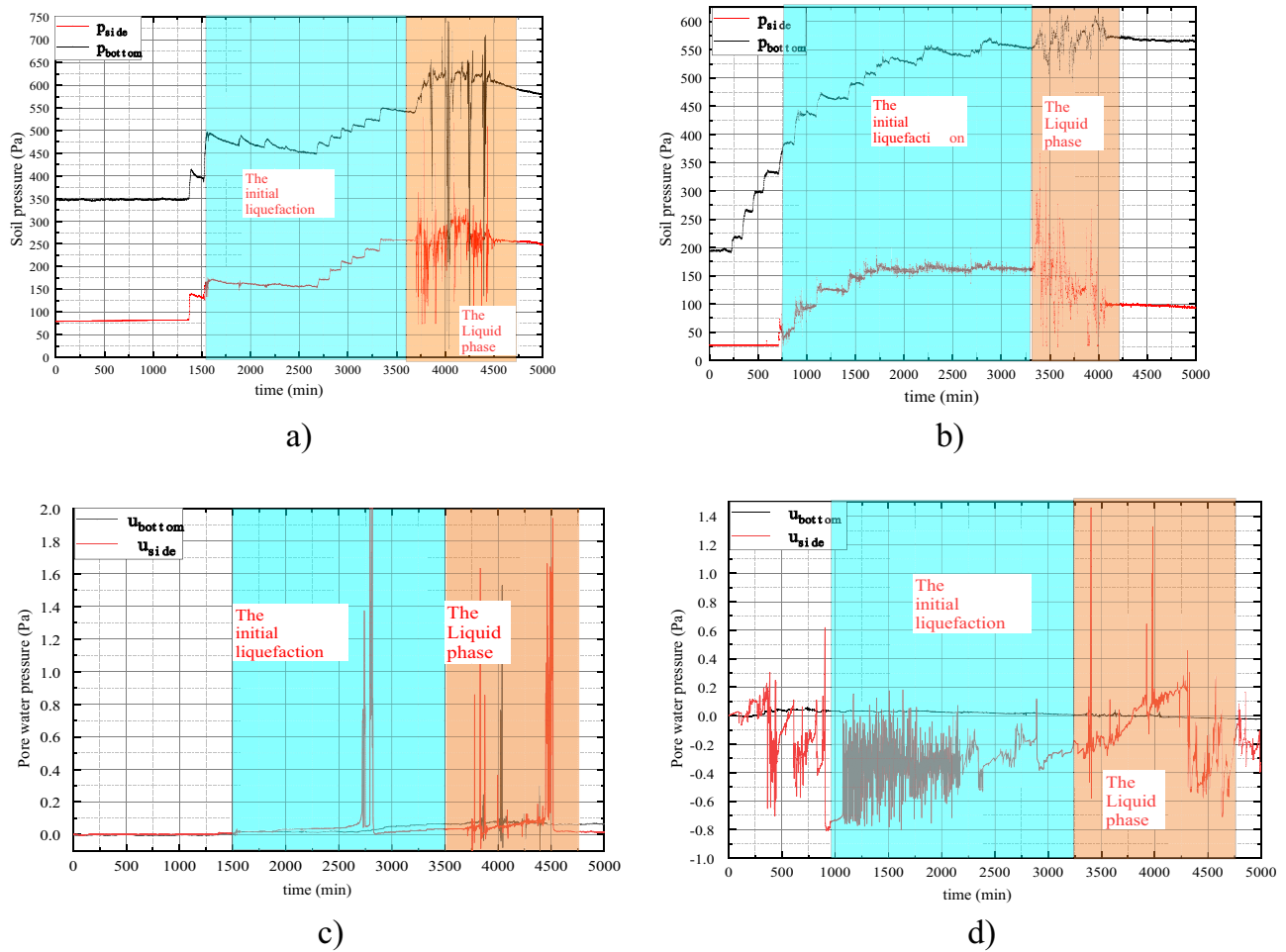
In this study, the occurrence of liquefaction was identified through a combined analysis of pore pressure ratio and visible deformation patterns obtained via PIV [16]. The key indicator for pore pressure was the excess pore pressure ratio, defined as

$$r_u = \frac{u}{\sigma'_0}$$

where  $u$  is the excess pore water pressure  $u$  and  $\sigma'_0$  is the initial vertical effective stress. Following common practice in dynamic centrifuge and shaking table experiments, liquefaction was considered to occur when  $r_u \geq 1.0$ , indicating a full loss of effective stress.

To complement the pore pressure analysis, PIV was used to observe deformation patterns such as sudden lateral spreading, flow-type displacements, or settlement





**Figure 4:** Sensor records under different disturbance conditions: (a) changes in soil pressure under the hammer condition, (b) changes in soil pressure under the rotating wheel condition, (c) changes in pore water pressure under the hammer condition, and (d) changes in pore water pressure under the rotating wheel condition.

acceleration. Specifically, we identified the onset of liquefaction when  $r_u \geq 0.95$  was sustained for more than two consecutive time steps.

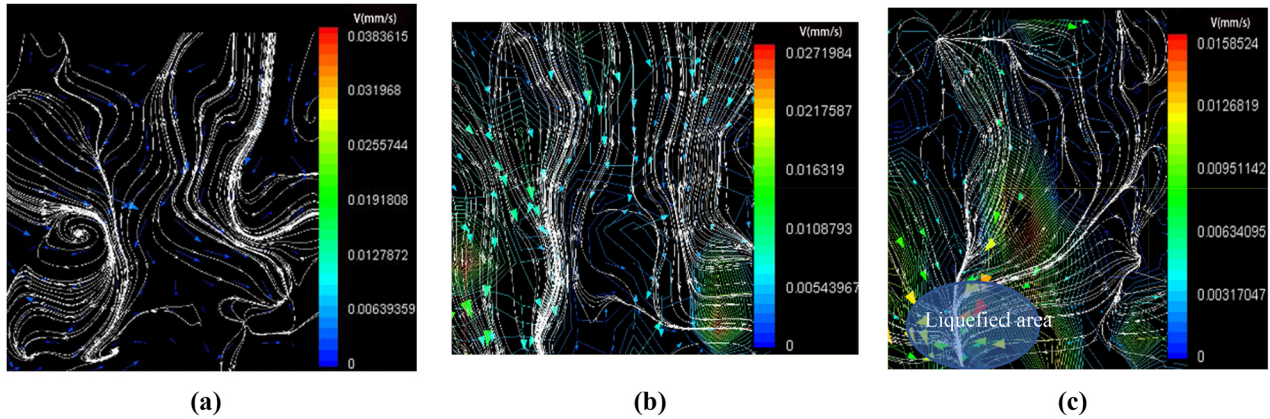
Simultaneously, PIV revealed a rapid increase in displacement rate and a loss of inter-particle constraint (e.g., sliding or shearing bands with strain concentration  $>5\%$ ).

The synchronization of both indicators ensures that transient pore pressure spikes without corresponding deformation were not misclassified as liquefaction. This dual-criteria approach enhances the robustness of liquefaction detection under complex boundary conditions.

Figure 4 displays sensor-recorded responses of soil and pore water pressure under two different disturbance conditions: hammer and rotating wheel. Subfigures (a) and (b) show the variation in soil pressure, while (c) and (d) illustrate changes in pore water pressure. The data indicate that rotating wheel disturbance produces more gradual and sustained pressure responses, suggesting deeper particle

rearrangement and pore water dissipation. These results are critical for understanding the mechanical behavior of silty soil under physical disturbance. The figure captions provide standalone interpretability to support scientific rigor.

According to the soil pressure sensor data,  $P$  shows a progressive increase at the initial hammering stage. The soil pressure also gradually increases with hammering at this stage due to dense particles. The pore water pressure  $u$  increases rapidly under the hammer condition with an increase rate of 100%. In particular, the lateral pore water pressure increases slowly from the initial stage to the late stage. When  $P$  reaches about 80%, the particle pores are squeezed to the limit at two hammering per minute, leading to a significant increase in  $u$ . The  $u_{\text{side}}$  and  $P_{\text{side}}$  are most significantly affected by hammering. After a certain accumulation of the hammering energy, the  $u$  increases rapidly, indicating soil liquefaction. As shown in Figure 4a and c, the abrupt change in  $u$  occurs earlier than the abrupt



**Figure 5:** PIV experiment of the flow field structure under the hammer condition: (a) start hammering, (b) soil layer before liquefaction, and (c) soil layer after liquefaction.

change in  $P$ . This phenomenon indicates that the dissipation direction of pore water pressure in the silty soil layer is vertically upward under hammer action. The intense pore water pressure causes dramatic fluctuations in soil pressure, making particles susceptible to re-suspension. Furthermore, concentrated and rapidly increasing pore pressure can lead to the rupture of the soil skeleton, followed by the re-suspension and re-deposition of coarse and fine particles.

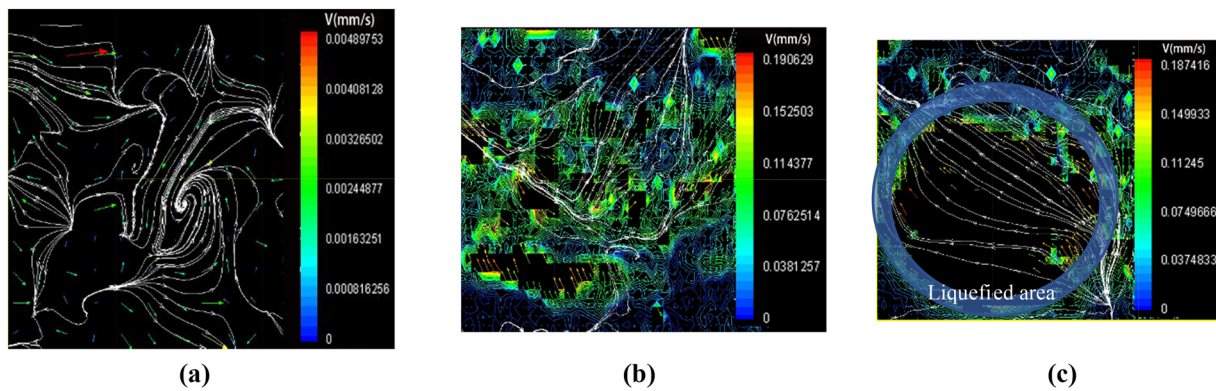
The disturbance data of the rotating wheel show that the  $P_{\text{side}}$  fluctuates significantly with the rotation, with a variation amplitude of 0.8 Pa. Moreover, the influence on the  $u_{\text{bottom}}$  by the rotating wheel is less due to the thick soil layer, with a float of only 0.1 Pa. According to the soil pressure analysis of the rotating wheel, the increase in the  $P_{\text{bottom}}$  relative to the  $P_{\text{side}}$  reaches 450 Pa after disturbance, higher than the increase under the hammer condition. When soil liquefaction occurs, the fluctuations generated by the rotating wheel are mainly in the liquefaction center, and the degree of bottom liquefaction is lower than that of hammer liquefaction. This result indicates that the degree and extent of hammer liquefaction are higher than those of rotating wheel liquefaction. As shown in Figure 4b and d, the significant change in soil pressure occurs earlier than that in pore pressure. This phenomenon suggests that wheel disturbance causes soil liquefaction by gradually increasing the soil pressure, resulting in the weakest fine soil particles vibrating out of the soil skeleton. After liquefaction, coarse soil particles are squeezed downward, promoting the upward movement of fine particles and significantly increasing the lateral pore pressure. Fine soil particles are compressed by coarse particles, creating upward pore water flow and migrating upward, further exacerbating the fluctuations in lateral soil pressure.

### 3.2 Field test and analysis of PIV

The PIV test system was used to analyze the velocity gradient (represented by velocity blocks in different colors) and the migration direction (represented by white arrows) of the particles. According to the analysis of the disturbance liquefaction by PIV, the hammer flow field moves up and down in the initial disturbance stage, while the rotating flow is swirling with most particles in a static state. The silt flow fields before and after liquefaction were recorded in conjunction with force field experiments.

Figures 5 and 6 present the PIV-observed flow field evolution of silty soil under two disturbance conditions: hammering and rotating wheel. Each figure captures the flow field at three stages – initial disturbance (a), pre-liquefaction (b), and post-liquefaction (c). Under hammering, the flow field remains localized with limited particle migration, while the rotating wheel generates a more uniform and extensive rearrangement of soil particles. These figures highlight the comparative effectiveness of disturbance types in promoting soil liquefaction and consolidation, with captions designed for standalone interpretation and enhanced scientific clarity.

A comparative analysis of the PIV-derived flow fields during the initial liquefaction stage (Figures 5 and 6) reveals distinct differences between the hammer-induced and rotating wheel-induced disturbances. Under hammering disturbance, a stratified velocity field develops, with a step-change in velocity between the lower and upper blocks. The lower layer exhibits a higher velocity of approximately 1.9 mm/s, which accelerates the downward migration and deposition of particles. In contrast, the rotating wheel disturbance generates flow primarily in the middle to upper soil layers, producing numerous small-



**Figure 6:** PIV experiment of the flow field structure under the rotating wheel condition: (a) start hammering, (b) soil layer before liquefaction, and (c) soil layer after liquefaction.

scale vortices. These vortices aggregate into vortex groups, with a maximum observed velocity of approximately 1.5 mm/s, facilitating turbulent mixing and moderate upward suspension of particles.

In the late stage of liquefaction, maximum particle velocity ( $v_{\max}$ ) under hammering increases to 2.2 mm/s, accompanied by high-amplitude particle fluctuations and the formation of clear migration channels throughout the liquefied layer. Particle streamlines show substantial vertical movement, with both upward and downward migration, contributing to secondary suspension and redeposition cycles. Conversely, in the rotating wheel scenario, the particle flow is characterized by projectile migration near the upper disturbance region and the formation of a projectile zone at the lower part of the vortex. Fine particles tend to be carried along the disturbance path, while coarse particles remain near their original positions, leading to enhanced particle separation and stratification.

Quantitative comparisons further support these observations. The hammering disturbance caused a larger increase in pore water pressure ( $\Delta u \approx 15.3$  kPa) compared to the rotating wheel disturbance ( $\Delta u \approx 9.7$  kPa). Additionally,  $d_{50}$  measurements after liquefaction indicated a more significant particle rearrangement under the hammering scenario (from 0.028 to 0.033 mm), suggesting downward segregation of coarser particles. In contrast, the rotating wheel test only showed a marginal change in  $d_{50}$  (from 0.028 to 0.030 mm), indicative of a finer-scale redistribution. Estimated shear stress ( $\tau$ ) during hammering reached peak values of 3.2 kPa, compared to 2.1 kPa under wheel rotation.

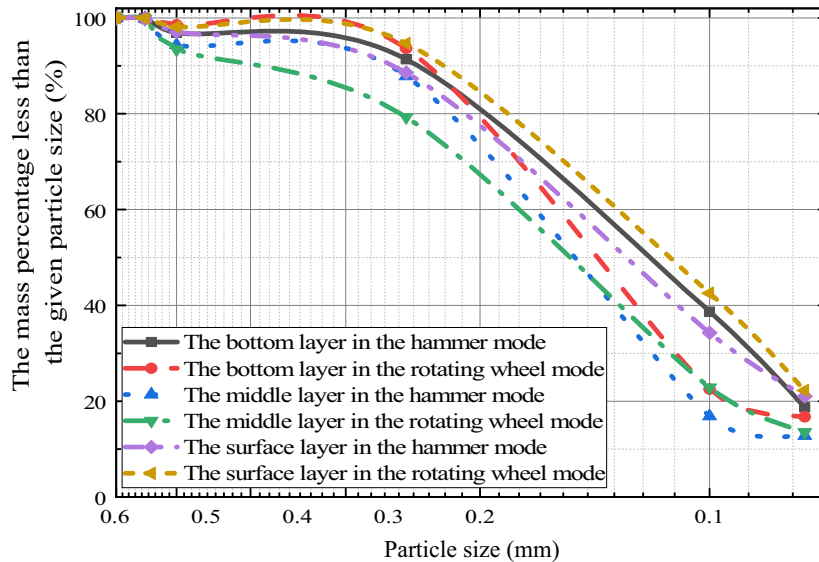
These findings indicate that hammering results in a higher energy input, promoting more vigorous particle migration and deeper soil reorganization, which may be beneficial for enhancing permeability but could delay structural stabilization due to repeated suspension. In

contrast, the rotating wheel produces localized disturbance and efficient particle separation, making it more suitable for controlled soil improvement where selective stratification and strength layering are desired.

### 3.3 Analysis of grading test results

The classification of soil samples in the late stages of the PIV test under different disturbance conditions is shown in Figure 7. The particles migrate and separate under the action of the hammer and rotating wheel, resulting in the re-migration of bottom particles with the pore water of the liquefied soil. In this experiment, the laser particle size analysis was conducted on the 0.1 m (surface layer), 0.3 m (middle layer), and 0.5 m (bottom layer) soil layers. The surface, middle, and bottom layers were disturbed by the rotating wheel, with median  $d_{50}$  of 0.12, 0.18, and 0.14 mm, respectively. Under the hammering disturbance, the medians  $d_{50}$  of the surface layer, middle layer, and bottom layer are 0.14, 0.15, and 0.12 mm, respectively. Compared to the bottom layer, the middle layer shows an increase in  $d_{50}$  particle size of 0.04 mm, and the surface layer exhibits a decrease in median  $d_{50}$  of nearly 0.05 mm. The highest clay content appears in the surface layer disturbed by the rotating wheel, accounting for 12% of particles. The fine particles are enriched, and particles in the middle layer gradually become coarser. However, the separation caused by hammering is characterized by a greater concentration of particle sizes. Despite the separation, the  $C_c$  values for the surface layer, middle, and bottom layers are 1.142, 1.512, and 1.422, respectively, indicating that the separation caused by hammering is insignificant. The reason is the excessive disturbing force of hammering, leading to the secondary lifting and deposition of particles. The  $C_c$  values for the surface, middle, and





**Figure 7:** Particle distribution curves of the soil sample.

bottom layers disturbed by the rotating wheel are 2.41, 1.371, and 1.333, respectively. This result suggests that the rotating wheel disturbance exacerbates the destruction of the soil skeleton, with fine particles gradually moving toward the surface due to the settlement of coarse particles and the upward flow induced by the disturbance force. As a result, the particle coarsening effect in the disturbance area is better.

### 3.4 Experiment on basic physical properties of the 1.2 m high flume

The model was extended to a 1.2 m-high flume bin, and the deposited stratum in each bin was left for 24 h under different disturbance liquefaction measures. The bearing capacity stratum of the natural sedimentary soil layer is 20 cm from the bottom, as measured by a portable bearing capacity tester. Therefore, the disturbed and hammering surfaces are the strata of this layer. Moreover, the liquefaction of the disturbed stratum was analyzed. The bearing capacity and its position and size were measured at 24 and 48 h after setting the disturbance static and dynamic forces. The initial strength was the reading of the load-bearing capacity tester. After 48 h, the corresponding soils were extracted to analyze the water content and density. Physical indicators were evaluated for the upper soft soil, hard-shell, and lower soft soil layers.

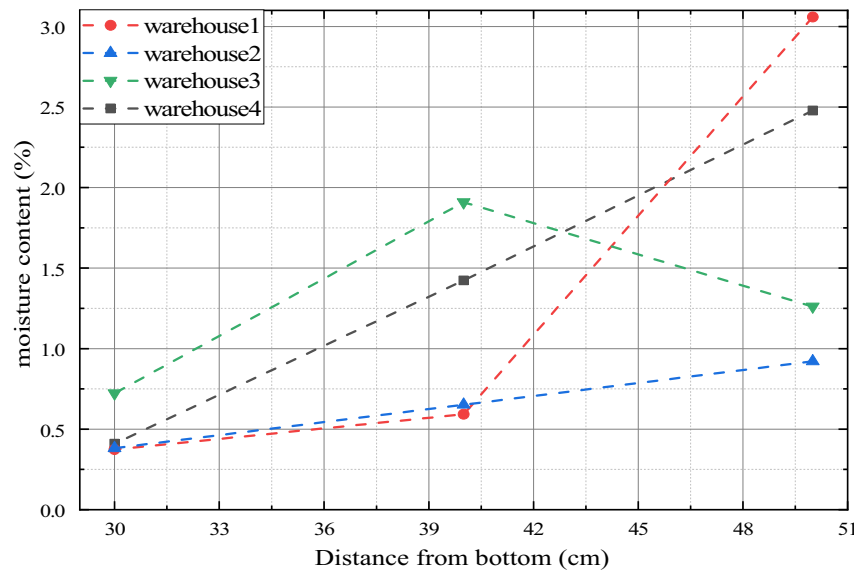
The test results of the water content and density of different layers and the bearing capacity of the hard-shell foundation formed after liquefaction are shown in Figures 8–11.

The hard-shell layers of Warehouses 1–4 are 5, 3, 10, and 0 cm, respectively. The initial strength and bearing capacity of the hard-shell layer in Warehouse 2 formed by hammer liquefaction are shown in Figure 8, with a maximum bearing capacity of 3.8 kN. Compared with Warehouse 4, Warehouse 2 failed to develop an initial strength within 12 h due to natural deposition, with an initial strength of only 0.8 at 48 h. In contrast, a thick shell layer is formed under the rotating wheel condition, with a load-bearing strength  $F$  of 2.8 kN and a maximum disturbance density  $\rho$  of  $2.9 \text{ g/mm}^3$ .

The above results also indicate that the physical properties of the hard-shell layers formed by different disturbance modes differ significantly. The high strength of the hard-shell layer formed by hammering disturbance is due to the compaction of the strata by hammering. However, its large disturbance area leads to particle re-suspension, resulting in a relatively thin hard-shell layer. In contrast, waves generated by rotation under the rotating wheel disturbance have a limited impact area. Wave disturbance leads to the migration of fine particles and the deposition of coarse particles. Due to the wide range of pore water pressure transmission, the thicker hard-shell layer has a relatively lower water content than the 50 cm surface layer.

### 3.5 Rheological test analysis

To investigate the rheological behavior of disturbed soil layers, a series of controlled rheological tests were conducted on siphoned samples from the upper, middle, and lower layers of disturbed soil columns. The soil samples



**Figure 8:** Variation of the water content.

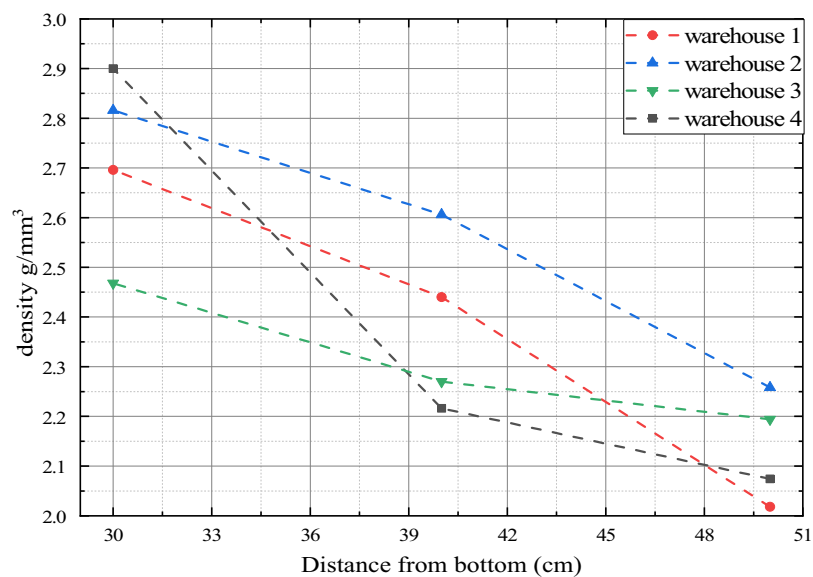
were extracted using a siphon device immediately following disturbance events, with minimal handling to preserve the *in situ* water content and microstructure.

Rheological testing was performed using a rotational rheometer (Model: Anton Paar MCR 302) equipped with a parallel plate measuring system (PP25). The test temperature was maintained at  $25 \pm 0.5^\circ\text{C}$ . Each soil sample was loaded into the rheometer with a gap setting of 1 mm, and pre-shearing was avoided to maintain the native structure of the disturbed material. The test involved a shear rate sweep from  $0.01$  to  $100\text{ s}^{-1}$ , and both stress-controlled

and strain-controlled modes were used to examine the yield stress, apparent viscosity, and viscoelastic deformation behavior.

Figure 12 presents typical stress–strain response curves from the tested soil samples. The analysis shows that the soil from bin 4 exhibited extremely low resistance under shear, indicative of high fluidity and weak inter-particle bonding – resulting in a near-zero recorded shear stress.

In the upper soil layer, disturbance led to reorganization and partial agglomeration of fine and coarse particles. The sample subjected to rotating wheel disturbance



**Figure 9:** Variation of density.

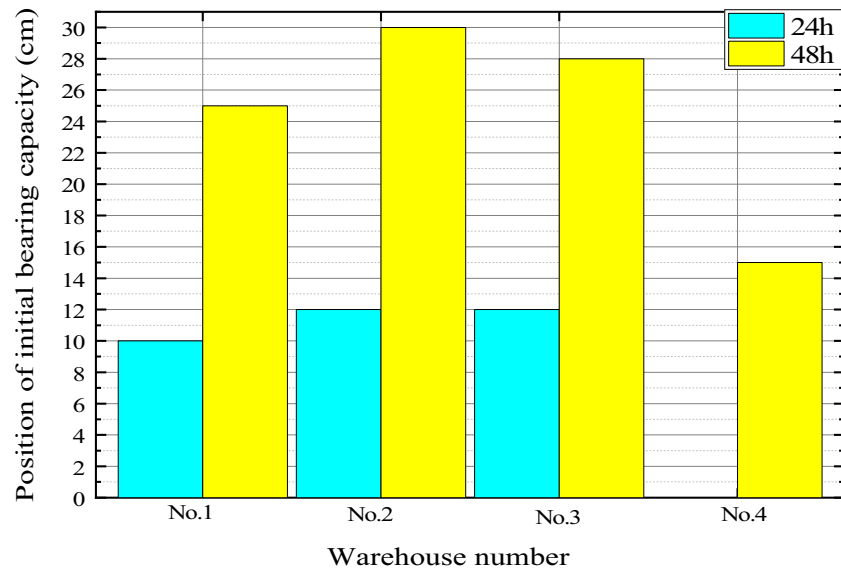


Figure 10: Development law of bearing capacity.

showed a maximum stress of 70 Pa, whereas hammering led to surface strength values around 10 Pa, significantly lower than that of undisturbed topsoil.

For the middle layer, the rotating wheel-induced hard-shell zone demonstrated the highest peak shear stress (approximately 40–50 Pa) and residual stiffness, indicating that micro-dynamic loading compacted soil particles effectively. In contrast, hammering-induced soil had lower residual strength due to the generation of high excess pore water pressure, which impeded structural reformation during the consolidation phase.

At the bottom layer, the soil beneath the hard-shell exhibited strength similar to that of the upper layer. However, under hammering disturbance, this layer showed approximately 50% higher residual strength than the upper soil, indicating that the higher-energy compaction from hammering led to significant soil densification.

Overall, rheological indices such as yield stress and apparent viscosity confirmed that rotational wheel disturbance led to better-structured hard-shell formation in the middle layer, while hammering caused deeper compaction

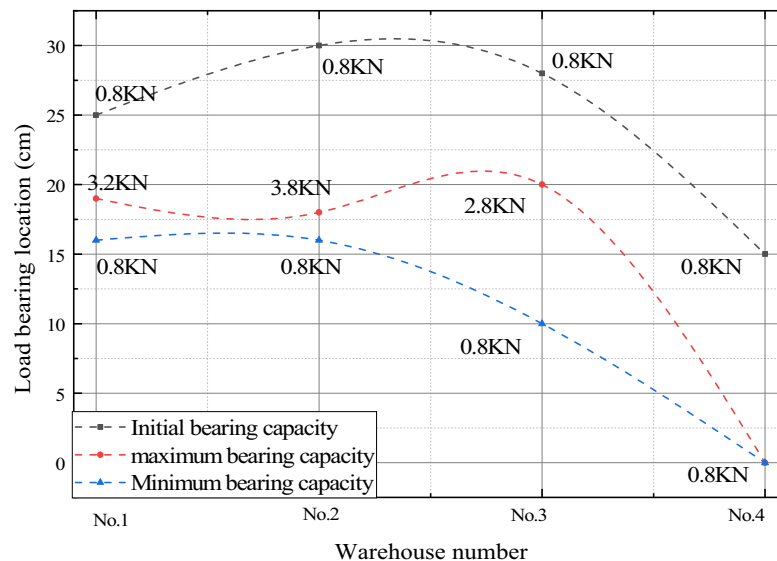
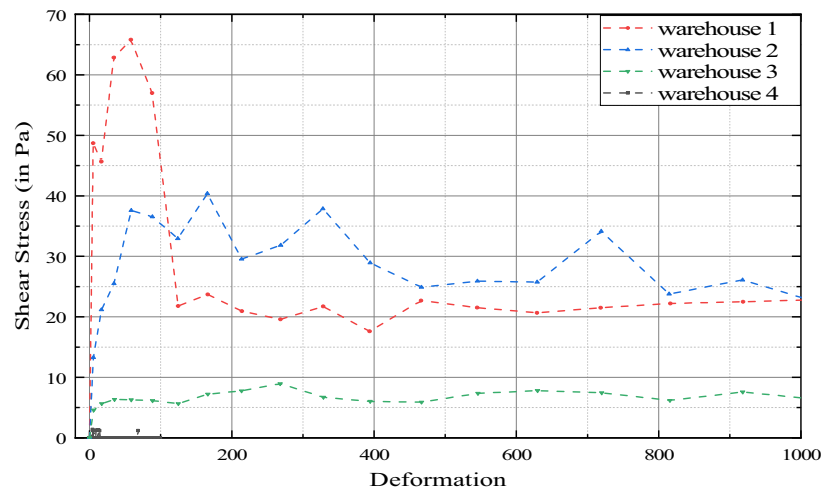
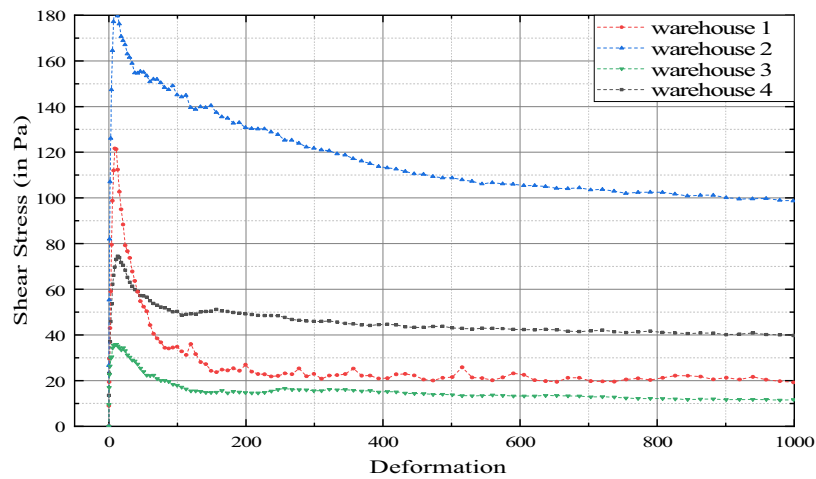


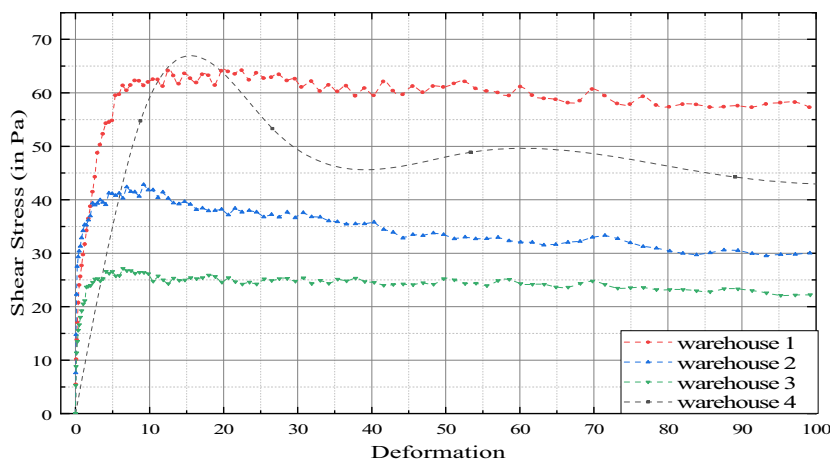
Figure 11: Distribution of bearing capacity.



a)



b)



c)

**Figure 12:** Rheological test results: (a) test results of the topsoil, (b) test results of the middle soil, and (c) test results of the subsoil change.



with stronger underlying soil layers. These findings are consistent with previous cyclic disturbance studies.

For future work, the rheological response of disturbed soils could be further explored using cyclic triaxial compression tests. Their approach offers valuable insight into the accumulation of pore pressure and strain under repetitive dynamic loading and may provide a more comprehensive framework for understanding long-term deformation behavior in disturbed subsoils.

Figure 12 presents the rheological test results of silty soil samples collected from different depths: (a) topsoil, (b) middle soil, and (c) subsoil. The curves reflect the variation in shear stress and viscosity under applied strain, indicating differences in structural strength and consolidation behavior across layers. The topsoil exhibits lower resistance, while deeper layers show higher stiffness and yield strength. These results demonstrate how depth influences the mechanical response of disturbed silty soils. The standardized captions ensure the figure is independently interpretable and aligned with scientific reporting standards.

## 4 Discussion

To evaluate the effectiveness of different artificial disturbance modes in promoting silty soil liquefaction and the subsequent formation of hard-shell layers, it is essential to ground the discussion in the classical mechanical framework of soil liquefaction [1]. Under cyclic loading, the contact network between saturated silt particles begins to break down. Two key factors – pore water pressure and soil pressure – govern the evolution of inter-particle forces and soil structure [17]. In particular, clay particles play a dual role: they act either as fillers that occupy void spaces or as agents of inter-particle cementation, depending on their location within the soil skeleton and the stage of disturbance.

### 4.1 Mechanism of hard-shell layer formation

During the initial phase of liquefaction, silt particles are compressed under cyclic stress, and fine clay particles are gradually squeezed into the interstitial spaces of the silt matrix. This transition changes their role from passive fillers to active cementing agents [18,19]. As pore water pressure builds and approaches the effective stress threshold, both silt and clay particles become suspended, losing contact and participating in random particle migration. Once this suspended state is reached, the dissipation of excess pore pressure allows for particle rearrangement, often

resulting in a denser and more stable configuration. It is in this post-liquefaction recovery phase that a hard-shell layer – defined as a zone with greater strength than the surrounding undisturbed soil – can emerge due to localized particle densification and new contact reformation.

## 4.2 Disturbance mode-specific behavior

### 4.2.1 Hammering disturbance

Experimental data (Figures 13 and 14) show that hammering disturbance rapidly increases both pore water pressure and soil pressure below the impact zone. As the duration of hammering increases, the influence range of the induced flow field expands vertically, driven by upward and rebound forces. This strong upward flow remobilizes previously settled silty particles and promotes the formation of a low-density, high-water-content suspended layer. Upon dissipation of pore pressure, this layer redeposits, and clay particles – now embedded into the lower portion of the soil skeleton – strengthen the lower part of the liquefied layer through cementation.

However, the middle part of the profile experiences a loss of fine silt particles, leading to the formation of a thin hard-shell layer with relatively high strength due to vertical densification. The upper layer, conversely, remains loosely packed, with mixed clay and silt redeposition, reflecting reduced structural integrity.

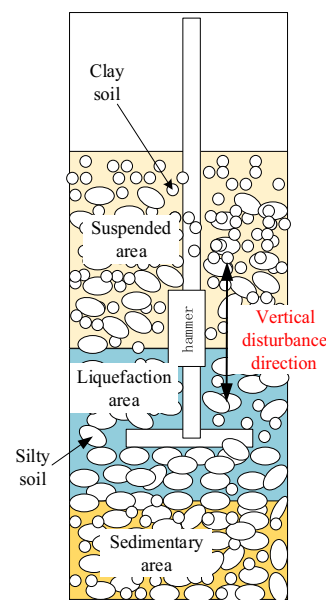


Figure 13: Schematic diagram of hammering disturbance.

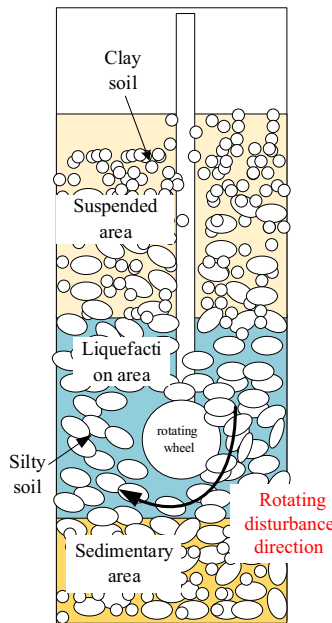


Figure 14: Schematic diagram of rotating wheel disturbance.

#### 4.2.2 Rotating wheel disturbance

Under the cyclic lateral loading caused by the rotating wheel, wave-like forces induce a periodic rise and fall of pore water pressure and lateral soil pressure. The lower soil layers tend to restrict vertical dissipation of pore pressure, which results in localized liquefaction near the rotating zone. In this case, small clay particles migrate upward through weakened inter-particle bonds, moving toward the upper soil layers [20]. As the liquefaction process completes, these particles precipitate into available pore spaces, rearranging the particle skeleton.

The result is a thick upper hard-shell layer that contains a high proportion of clay. However, due to the low density and high water content, this shell exhibits inferior mechanical strength compared to the thinner but denser layer formed under hammering. Meanwhile, the bottom of this hard-shell layer remains relatively unaffected, while the soil immediately beneath the rotating wheel exhibits high density and improved strength, benefitting from cyclic compaction and reduced pore space.

### 4.3 Comparative mechanistic insight

The comparative mechanisms can be summarized as follows:

- (1) Hammering produces deep, vertical particle migration, promoting a denser and thinner hard-shell layer in the

middle zone. The short migration path and high-energy input contribute to rapid consolidation and local strength gain.

- (2) Rotating wheel induces horizontal, surface-proximal particle migration, leading to a thicker, clay-rich upper layer with lower mechanical strength, but more uniform compaction in the sub-layer.

Thus, while both methods can induce the formation of hard-shell structures, the hammering method is more effective for targeted deep strengthening, whereas the rotating wheel is suitable for broader, surface-level soil improvement.

## 5 Conclusion

This study investigated the formation mechanism of hard-shell layers in saturated silty soils subjected to two distinct artificial disturbance modes: vertical hammering and rotating wheel-induced cyclic shear. The experiments aimed to simulate post-liquefaction recovery processes relevant to estuarine and coastal environments affected by wave and seismic activity. The key conclusions are as follows.

Hard-shell layer formation is inherently linked to the fluid – plastic transition of saturated silty soils during liquefaction. In this state, vibration and disturbance allow fine clay particles to migrate into the pore spaces among silt grains. These particles act primarily as mechanical fillers rather than cementing agents, resulting in a densely packed soil skeleton dominated by silt. Upon pore pressure dissipation, this structure re-stabilizes into a rigid layer with increased density and shear resistance.

Comparative analysis reveals distinct differences in liquefaction behavior and post-liquefaction structure between hammering and rotating wheel disturbances. Hammering induces a rapid and substantial increase in pore water pressure – doubling initial values – and achieves a high degree of liquefaction. However, the high-energy impact also causes strong particle re-suspension, which may prevent complete precipitation of clay particles. This results in a thin hard-shell layer (~2 cm) with high compaction in the underlying deposit zone. In contrast, rotating wheel disturbance produces moderate liquefaction but forms more stable vortex flow fields that facilitate selective particle separation. This leads to a thicker hard-shell layer (~10 cm), albeit with slightly lower strength. Notably, the compression index ( $C_c$ ) of liquefied soil under wheel disturbance was

reduced to 1.33, suggesting improved consolidation behavior.

There is a clear trade-off between shell thickness and mechanical strength. Hammering yields a thin but highly compacted shell with a penetration strength of up to 3.6 N, offering high bearing capacity in localized zones. Meanwhile, rotating wheel disturbance forms a thicker shell that may provide broader coverage but has lower strength due to higher water content and less uniform particle packing. This trade-off suggests that hammering is more suitable for deep reinforcement applications, while rotating wheel methods may be advantageous for improving surface stability over larger areas.

Overall, this study provides mechanistic insight into liquefaction-driven soil reorganization and offers practical guidance for tailoring artificial disturbance methods to achieve targeted subsurface strengthening in coastal silty environments.

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participated in the discussion of research results; proofread the final version of the manuscript.

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**Data availability statement:** The data used to support the findings of this study are included in the article.

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